

Trigonometric and elliptic Ruijsenaars–Schneider systems on the complex projective space

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Abstract

We present a direct construction of compact real forms of the trigonometric and elliptic n -particle Ruijsenaars–Schneider systems whose completed center-of-mass phase space is the complex projective space \mathbb{CP}^{n-1} with the Fubini–Study symplectic structure. These systems are labelled by an integer $p \in \{1, \dots, n-1\}$ relative prime to n and a coupling parameter y varying in a certain punctured interval around $p\pi/n$. Our work extends Ruijsenaars’s pioneering study of compactifications that imposed the restriction $0 < y < \pi/n$, and also builds on an earlier derivation of more general compact trigonometric systems by Hamiltonian reduction.

1 Introduction

The investigation of integrable systems of particles moving in one spatial dimension started decades ago and persistently attracts intense attention due to the fascinating mathematics and diverse physical applications of these systems, as reviewed in [4, 10, 11, 15, 17, 20]. The Ruijsenaars–Schneider (RS) model [13, 14] occupies a central position in this family, since many other interesting models of Calogero–Moser–Sutherland and Toda type can be obtained from it as various limits and analytic continuations [17]. The phase space of these particle systems is usually the cotangent bundle of the configuration space, which is never compact due to the infinite range of the canonical momenta. The standard RS Hamiltonian depends on the momenta ϕ_k through the function $\cosh(\phi_k)$, but by analytic continuation this may be replaced by $\cos(\phi_k)$, which effectively compactifies the momenta on a circle. If the dependence on the position variables x_k is also through a periodic function, then the phase space can be taken to be a bounded set. This possibility was examined in [16], where the Hamiltonian

$$H(x, \phi) = \sum_{k=1}^n \cos(\phi_k) \sqrt{\prod_{\substack{j=1 \\ (j \neq k)}}^n \left[1 - \frac{\sin^2 y}{\sin^2(x_j - x_k)} \right]} \quad (1.1)$$

containing a real coupling parameter $0 < y < \pi/2$ was considered. Ruijsenaars called this the III_b system, with III referring to the trigonometric character of the interaction, as in [11], and the suffix standing for ‘bounded’. (One may also introduce another real parameter into the III_b system, by replacing ϕ_k say by $\beta\phi_k$.) The domain of the ‘angular position variables’ $\{(x_1, \dots, x_n)\} \subset [0, \pi]^n$ must be restricted in such a way that the Hamiltonian (1.1) is real and smooth. This may be ensured by prescribing

$$x_{i+1} - x_i > y \quad (i = 1, \dots, n-1), \quad x_n - x_1 < \pi - y, \quad (1.2)$$

which obviously implies Ruijsenaars’s condition

$$0 < y < \frac{\pi}{n}. \quad (1.3)$$

Although the Hamiltonian is then real, its flow is not complete on the naive phase space, because it may reach the boundary $x_{k+1} - x_k = y$ (with $x_{k+n} \equiv x_k + \pi$) at finite time [16]. Completeness of the commuting flows is a crucial property of any bona fide integrable system, but one cannot directly add the boundary to the phase space because that would not yield a smooth manifold. One of the seminal results of [16] is the solution of this conundrum. In fact, Ruijsenaars constructed a symplectic embedding of the center-of-mass phase space of the system into the complex projective space \mathbb{CP}^{n-1} , such that the image of the embedding is a dense open submanifold and the Hamiltonian (1.1) as well as its commuting family extend to smooth functions on the full \mathbb{CP}^{n-1} . As \mathbb{CP}^{n-1} is compact, the corresponding Hamiltonian flows are complete. The resulting ‘compactified trigonometric RS system’ has been studied at the classical level in detail [16], and after an initial exploration of the rank 1 case [15], its quantum mechanical version was also solved [19]. These classical systems are self-dual in the sense that their position and action variables can be exchanged by a canonical transformation of order 4, somewhat akin to the mapping $(x, \phi) \mapsto (-\phi, x)$ for a free particle, and their quantum mechanical versions enjoy the bispectral property [15, 19].

The possibility of an analogous compactification of the elliptic RS system having the Hamiltonian

$$H(x, \phi) = \sum_{k=1}^n \cos(\phi_k) \sqrt{\prod_{\substack{j=1 \\ (j \neq k)}}^n [s(y)^2 (\wp(y) - \wp(x_j - x_k))]} \quad (1.4)$$

with functions \wp (4.1) and s (4.2) was pointed out in [15, 17], but it was not described in detail.

Even though it was only proved [16] that the restrictions (1.2), (1.3) are sufficient to allow compactification, equation (1.3) was customarily mentioned in the literature [6, 9, 15, 17, 19] as a necessary condition for the systems to make sense. However, in a recent work [8] a completion of the III_b system on a compact phase space was obtained for any generic parameter

$$0 < y < \pi. \quad (1.5)$$

The paper [8] relied on deriving compactified RS systems in the center-of-mass frame via reduction of a ‘free system’ on the quasi-Hamiltonian [1] double $\text{SU}(n) \times \text{SU}(n)$. This was achieved by setting the relevant group-valued moment map equal to the constant matrix $\mu_0(y) = \text{diag}(e^{2iy}, \dots, e^{2iy}, e^{-2(n-1)iy})$, and it makes perfect sense for any (generic) y . The corresponding domain of the position variables depends on y and differs from the one posited in (1.2). The possibility to relax the condition (1.3) on y also appeared in [3].

The principal motivation for our present work comes from the classification of the coupling parameter y found in [8]. Namely, it turned out that the reduction is applicable except for a finite set of y -values, and the rest of the set $(0, \pi)$ decomposes into two subsets, containing so-called type (i) and type (ii) y -values. The ‘main reduced Hamiltonian’ always takes the III_b form (1.1) on a dense open subset of the reduced phase space. In the type (i) cases the particles cannot collide and the action variables of the reduced system naturally engender an isomorphism with the Hamiltonian toric manifold \mathbb{CP}^{n-1} . In type (ii) cases, that exist for any $n > 3$, the reduction constraints admit solutions $(a, b) \in \text{SU}(n) \times \text{SU}(n)$ for which the eigenvalues of a or b are not all distinct, entailing that the particles of the reduced system can collide. For a detailed exposition of these succinct statements, the reader may consult [8]. We here only add the remark that the connected domain of the positions always contains the equal-distance configuration $x_{k+1} - x_k = \pi/n$ ($\forall k$) for which the number of negative factors in each product under the square root in (1.1) is $2\lfloor ny/\pi \rfloor$ if $0 < y < \pi/2$ and $2\lfloor n(\pi - y)/\pi \rfloor$ if $\pi/2 < y < \pi$.

This Letter is exclusively concerned with the type (i) cases just mentioned. Our first goal is to reconstruct the corresponding compactification on \mathbb{CP}^{n-1} using only direct, elementary methods, i.e., not relying on reduction techniques. Such construction was not known previously except for the special type (i) cases (1.3), which we shall generalize. By doing so, we shall gain a better understanding of the structure of these trigonometric systems. This part of the Letter fills Sections 2 and 3 that follow. In Section 4, we explain that the direct method is applicable to obtain type (i) compactifications of the elliptic RS system as well. This new result extends the remarks of Ruijsenaars [15, 17].

It would have been possible to organize our text differently, starting with the elliptic case and then recovering the trigonometric systems as a limit. We opted for first presenting the trigonometric case for the reason that in our hope this makes the paper easier to understand, and also since this actually follows our line of research.

Our results lead to several open questions and possible topics for future work that will be outlined at the end of the paper.

2 Embedding of the local phase space into \mathbb{CP}^{n-1}

In this section we first recall the local phase space of the III_b model from [8], and then present its symplectic embedding into \mathbb{CP}^{n-1} in every type (i) case.

The III_b model can be thought of as n interacting particles on the unit circle with positions $\delta_k = e^{2ix_k}$. We impose the condition $\prod_{k=1}^n \delta_k = 1$, which means that we work in the ‘center-of-mass frame’, and parametrize the positions as

$$\delta_1(\xi) = e^{\frac{2i}{n} \sum_{j=1}^n j\xi_j}, \quad \delta_k(\xi) = e^{2i\xi_{k-1}} \delta_{k-1}(\xi), \quad k = 2, \dots, n, \quad (2.1)$$

where ξ belongs to a certain open subset \mathcal{A}_y^+ inside the ‘Weyl alcove’

$$\mathcal{A} = \{\xi \in \mathbb{R}^n \mid \xi_k \geq 0 \ (k = 1, \dots, n), \ \xi_1 + \dots + \xi_n = \pi\}. \quad (2.2)$$

Note that \mathcal{A} is a simplex in the $(n-1)$ -dimensional affine space

$$E = \{\xi \in \mathbb{R}^n \mid \xi_1 + \dots + \xi_n = \pi\}. \quad (2.3)$$

The local phase space can be described as the product manifold

$$P_y^{\text{loc}} = \{(\xi, e^{i\theta}) \mid \xi \in \mathcal{A}_y^+, \ e^{i\theta} \in \mathbb{T}^{n-1}\}, \quad (2.4)$$

where \mathbb{T}^{n-1} is the $(n-1)$ -torus, equipped with the standard symplectic form

$$\omega^{\text{loc}} = \sum_{k=1}^{n-1} d\theta_k \wedge d\xi_k. \quad (2.5)$$

The dynamics is governed by the Hamiltonian

$$H_y^{\text{loc}}(\xi, \theta) = \sum_{j=1}^n \cos(\theta_j - \theta_{j-1}) \sqrt{\prod_{m=j+1}^{j+n-1} \left[1 - \frac{\sin^2 y}{\sin^2(\sum_{k=j}^{m-1} \xi_k)} \right]}. \quad (2.6)$$

Here, $\theta_0 = \theta_n = 0$ have been introduced and the indices are understood modulo n , i.e.,

$$\xi_{m+n} = \xi_m, \quad \forall m. \quad (2.7)$$

The product under the square root is positive for every $\xi \in \mathcal{A}_y^+$, and thus $H_y^{\text{loc}} \in C^\infty(P_y^{\text{loc}})$. This model was considered in [8] for any y chosen from the interval $(0, \pi)$ except the excluded values that satisfy $e^{2imy} = 1$ for some $m = 1, \dots, n$.

According to [8], there are two different kinds of intervals for y to be in, named type (i) and (ii). The type (i) couplings can be described as follows. For a fixed positive integer $n \geq 2$, choose $p \in \{1, \dots, n-1\}$ to be a coprime to n , i.e., $\gcd(n, p) = 1$, and let q denote the multiplicative inverse of p in the ring \mathbb{Z}_n , that is $pq \equiv 1 \pmod{n}$. Then the parameter y can take its values according to either

$$\left(\frac{p}{n} - \frac{1}{nq}\right)\pi < y < \frac{p\pi}{n} \quad \text{or} \quad \frac{p\pi}{n} < y < \left(\frac{p}{n} + \frac{1}{(n-q)n}\right)\pi. \quad (2.8)$$

For such a type (i) parameter y , the local configuration space \mathcal{A}_y^+ is the interior of a simplex \mathcal{A}_y in E (2.3) bounded by the hyperplanes

$$\xi_j + \dots + \xi_{j+p-1} = y, \quad j = 1, \dots, n, \quad (2.9)$$

where (2.7) is understood. To give a more detailed description of \mathcal{A}_y , we introduce

$$M = p\pi - ny, \quad (2.10)$$

and note that (2.8) gives $M > 0$ and $M < 0$, respectively. Then any $\xi \in \mathcal{A}_y$ must satisfy

$$\text{sgn}(M)(\xi_j + \cdots + \xi_{j+p-1} - y) \geq 0, \quad j = 1, \dots, n. \quad (2.11)$$

In terms of the particle coordinates x_k , which are ordered as $x_{k+1} \geq x_k$ and extended by the convention $x_{k+n} = x_k + \pi$, the above condition says that

$$x_{j+p} - x_j \geq y \quad \text{if } M > 0 \quad \text{and} \quad x_{j+p} - x_j \leq y \quad \text{if } M < 0 \quad (2.12)$$

for every j . Therefore the distances of the p -th neighbouring particles on the circle are constrained. The n vertices of the simplex \mathcal{A}_y are explicitly given in [8] (Proposition 11 and Lemma 8 *op. cit.*). Every vertex and thus \mathcal{A}_y itself lies inside the larger simplex \mathcal{A} (2.2), entailing that $x_{j+1} - x_j$ possesses a positive lower bound in each type (i) case.

The type (ii) cases correspond to those admissible y -values that do not satisfy (2.8) for any p relative prime to n . In such cases \mathcal{A}_y^+ has a different structure [8]. Type (ii) cases exist for every $n \geq 4$. See Figure 1 for an illustration.

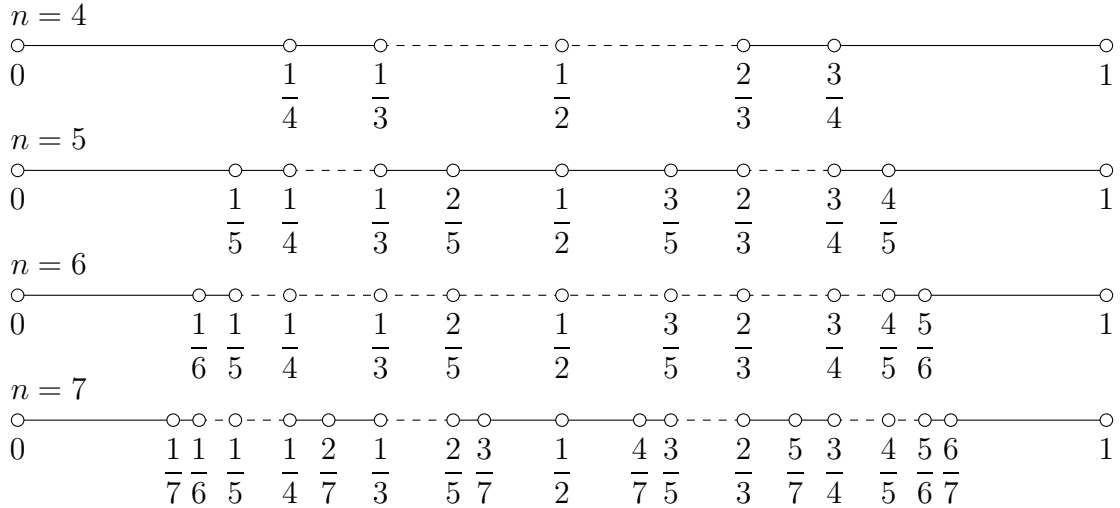


Figure 1: The range of y/π for $n = 4, 5, 6, 7$. The displayed numbers are excluded values. Admissible values of y form intervals of type (i) (solid) and type (ii) (dashed) couplings.

We further continue with the assumption that y satisfies (2.8). Motivated by [16, 6], we now introduce the map

$$\mathcal{E}: \mathcal{A}_y^+ \times \mathbb{T}^{n-1} \rightarrow \mathbb{C}^n, \quad (\xi, e^{i\theta}) \mapsto (u_1, \dots, u_n) \quad (2.13)$$

with the complex coordinates having the squared absolute values

$$|u_j|^2 = \text{sgn}(M)(\xi_j + \cdots + \xi_{j+p-1} - y), \quad j = 1, \dots, n, \quad (2.14)$$

and the arguments

$$\arg(u_j) = \text{sgn}(M) \sum_{k=1}^{n-1} \Omega_{j,k} \theta_k, \quad j = 1, \dots, n-1, \quad \arg(u_n) = 0, \quad (2.15)$$

where the $\Omega_{j,k}$ ($j, k = 1, \dots, n-1$) are integers chosen in such a way that

$$\mathcal{E}^* \left(i \sum_{j=1}^n d\bar{u}_j \wedge du_j \right) = \sum_{k=1}^{n-1} d\theta_k \wedge d\xi_k. \quad (2.16)$$

In order for (2.16) to be achieved Ω has to be the inverse transpose of the $(n-1) \times (n-1)$ coefficient matrix of ξ_1, \dots, ξ_{n-1} extracted from eqs. (2.14) by applying $\xi_1 + \dots + \xi_n = \pi$. In other words, the squared absolute values $|u_j|^2$ are written as

$$|u_j|^2 = \begin{cases} \operatorname{sgn}(M) \left(\sum_{k=1}^{n-1} A_{j,k} \xi_k - y \right), & \text{if } 1 \leq j \leq n-p, \\ \operatorname{sgn}(M) \left(\sum_{k=1}^{n-1} A_{j,k} \xi_k - y + \pi \right), & \text{if } n-p < j \leq n-1, \end{cases} \quad (2.17)$$

where A stands for the above-mentioned coefficient matrix, which has the components

$$A_{j,k} = \begin{cases} +1, & \text{if } 1 \leq j \leq n-p \text{ and } j \leq k < j+p, \\ -1, & \text{if } n-p < j \leq n-1 \text{ and } j+p-n \leq k < j, \\ 0, & \text{otherwise.} \end{cases} \quad (2.18)$$

A close inspection of the structure of A reveals that

$$\det(A) = (-1)^{(n-p)(p-1)} \prod_{j=1}^{n-p} A_{j,j+p-1} \prod_{k=1}^{p-1} A_{n-p+k,k} = (-1)^{(n-p+1)(p-1)} = +1, \quad (2.19)$$

therefore $\Omega = (A^{-1})^\top$ exists and consists of integers, as required in (2.15). Next, we give Ω explicitly.

Proposition 2.1. *The transpose of the inverse of the matrix A (2.18) can be written as*

$$\Omega = B - C, \quad (2.20)$$

where B is a $(0,1)$ -matrix of size $(n-1)$ with zeros along certain diagonals given by

$$B_{m,k} = \begin{cases} 0, & \text{if } k - m \equiv \ell p \pmod{n} \text{ for some } \ell \in \{1, \dots, n-q\}, \\ 1, & \text{otherwise,} \end{cases} \quad (2.21)$$

and C is also a binary matrix of size $(n-1)$ with zeros along columns given by

$$C_{m,k} = \begin{cases} 0, & \text{if } k \equiv \ell p \pmod{n} \text{ for some } \ell \in \{1, \dots, n-q\}, \\ 1, & \text{otherwise.} \end{cases} \quad (2.22)$$

Proof. We start by presenting a useful auxiliary statement. Let us introduce the subsets S and S_i of the ring \mathbb{Z}_n as

$$S = \{\ell p \pmod{n} \mid \ell = 1, \dots, n-q\}, \quad S_i = \{i + \ell \pmod{n} \mid \ell = 0, \dots, p-1\}, \quad (2.23)$$

for any $i \in \mathbb{Z}_n$. Then define $I_i \in \mathbb{N}$ to be the number of elements in the intersection $S_i \cap S$. Notice that $i \in S$ if and only if $(i+p) \in S$ except for $i \equiv (n-1) \equiv (n-q)p \pmod{n}$, for which $(n-1) + p \equiv (n-q+1)p \pmod{n}$ does not belong to S . It follows that

$$I_1 = \dots = I_{n-1} = I_n + 1. \quad (2.24)$$

Our aim is to show that $(A\Omega^\top)_{j,m} = \delta_{j,m}$ ($\forall j, m$) with Ω defined by (2.20)-(2.22). First, by the formula of A (2.18) for any $1 \leq j \leq n-p$ and $1 \leq m \leq n-1$ we have

$$(A\Omega^\top)_{j,m} = \sum_{k=1}^{n-1} A_{j,k} \Omega_{m,k} = \sum_{k=j}^{j+p-1} \Omega_{m,k} = \sum_{k=j}^{j+p-1} (B_{m,k} - C_{m,k}). \quad (2.25)$$

The definition of the matrices B (2.21) and C (2.22) gives directly that

$$\sum_{k=j}^{j+p-1} B_{m,k} = p - I_{j-m}, \quad \sum_{k=j}^{j+p-1} C_{m,k} = p - I_j. \quad (2.26)$$

By using (2.24), this readily implies that $(A\Omega^\top)_{j,m} = \delta_{j,m}$ holds for the case at hand.

Second, for any $n-p < j \leq n-1$ and $1 \leq m \leq n-1$ we have

$$(A\Omega^\top)_{j,m} = \sum_{k=1}^{n-1} A_{j,k} \Omega_{m,k} = \sum_{k=j+p-n}^{j-1} (-1) \Omega_{m,k} = \sum_{k=j+p-n}^{j-1} (C_{m,k} - B_{m,k}). \quad (2.27)$$

From this point on the reasoning is quite similar to the previous case, and we obtain that $(A\Omega^\top)_{j,m} = \delta_{j,m}$ always holds. \square

To enlighten the geometric meaning of the map \mathcal{E} (2.13), notice from (2.14) that

$$\sum_{j=1}^n |u_j|^2 = \text{sgn}(M)(p(\xi_1 + \dots + \xi_n) - ny) = \text{sgn}(M)(p\pi - ny) = |M|. \quad (2.28)$$

Then represent the complex projective space \mathbb{CP}^{n-1} as

$$\mathbb{CP}^{n-1} = S_{|M|}^{2n-1} / U(1) \quad (2.29)$$

with

$$S_{|M|}^{2n-1} = \{(u_1, \dots, u_n) \in \mathbb{C}^n \mid |u_1|^2 + \dots + |u_n|^2 = |M|\}. \quad (2.30)$$

Correspondingly, let

$$\pi_{|M|} : S_{|M|}^{2n-1} \rightarrow \mathbb{CP}^{n-1} \quad (2.31)$$

denote the natural projection and equip \mathbb{CP}^{n-1} with the rescaled Fubini–Study symplectic form $|M|\omega_{\text{FS}}$ characterized by the relation

$$\pi_{|M|}^*(|M|\omega_{\text{FS}}) = i \sum_{j=1}^n d\bar{u}_j \wedge du_j, \quad (2.32)$$

where the u_j 's are regarded as functions on $S_{|M|}^{2n-1}$. It is readily seen from the definitions that the map

$$\pi_{|M|} \circ \mathcal{E} : \mathcal{A}_y^+ \times \mathbb{T}^{n-1} \rightarrow \mathbb{CP}^{n-1} \quad (2.33)$$

is smooth, injective and its image is the open submanifold for which $\prod_{j=1}^n |u_j|^2 \neq 0$. Equations (2.5), (2.16) and (2.32) together imply the symplectic property

$$(\pi_{|M|} \circ \mathcal{E})^*(|M|\omega_{\text{FS}}) = \omega^{\text{loc}}, \quad (2.34)$$

from which it follows that this map is an *embedding*.

To summarize, in this section we have constructed the symplectic diffeomorphism $\pi_{|M|} \circ \mathcal{E}$ between the local phase space P_y^{loc} (2.4) and the dense open submanifold of \mathbb{CP}^{n-1} on which the product of the homogeneous coordinates is nowhere zero. If desired, the explicit formula of the smooth inverse mapping can be easily found as well.

3 Global extension of the trigonometric Lax matrix

It was proved in [8] with the aid of quasi-Hamiltonian reduction that the global phase space of the III_b model is \mathbb{CP}^{n-1} for the type (i) couplings, which we continue to consider. Here, we utilize the symplectic embedding (2.33) to construct a global Lax matrix on \mathbb{CP}^{n-1} explicitly, starting from the local RS Lax matrix defined on $\mathcal{A}_y^+ \times \mathbb{T}^{n-1}$. This issue was not investigated previously except for the $p = 1$ case of (2.8), see [16, 6, 8].

The local Lax matrix $L_y^{\text{loc}}(\xi, e^{i\theta}) \in \text{SU}(n)$ used in [8] contains the trigonometric Cauchy matrix C_y given with the help of (2.1) by

$$C_y(\xi)_{j,\ell} = \frac{e^{iy} - e^{-iy}}{e^{iy}\delta_j(\xi)^{1/2}\delta_\ell(\xi)^{-1/2} - e^{-iy}\delta_j(\xi)^{-1/2}\delta_\ell(\xi)^{1/2}}. \quad (3.1)$$

Thanks to the relation $\delta_k(\xi) = e^{2ix_k}$, this is equivalent to

$$C_y(\xi)_{j,\ell} = \frac{\sin(y)}{\sin(x_j - x_\ell + y)}. \quad (3.2)$$

Then we have

$$L_y^{\text{loc}}(\xi, e^{i\theta})_{j,\ell} = C_y(\xi)_{j,\ell} v_j(\xi, y) v_\ell(\xi, -y) \rho(\theta)_\ell, \quad \forall (\xi, e^{i\theta}) \in \mathcal{A}_y^+ \times \mathbb{T}^{n-1}, \quad (3.3)$$

where $\rho(\theta)_\ell = e^{i(\theta_{\ell-1} - \theta_\ell)}$ (applying $\theta_0 = \theta_n = 0$) and

$$v_\ell(\xi, \pm y) = \sqrt{z_\ell(\xi, \pm y)} \quad \text{with} \quad z_\ell(\xi, \pm y) = \text{sgn}(\sin(ny)) \prod_{m=\ell+1}^{\ell+n-1} \frac{\sin(\sum_{k=\ell}^{m-1} \xi_k \mp y)}{\sin(\sum_{k=\ell}^{m-1} \xi_k)}. \quad (3.4)$$

A key point [8] (which is detailed below) is that $z_\ell(\xi, \pm y)$ is positive for any $\xi \in \mathcal{A}_y^+$. We note for clarity that z_ℓ and v_ℓ above differ from those in [8] by a harmless multiplicative constant, and also mention that L_y^{loc} is a specialization of (a similarity transform of) the standard RS Lax matrix [17].

The spectral invariants of L_y^{loc} (3.3) yield a Poisson commuting family of functional dimension $(n-1)$ [17, 8], containing the Hamiltonian H_y^{loc} (2.6) due to the equation

$$\text{Re}(\text{tr} L_y^{\text{loc}}(\xi, e^{i\theta})) = H_y^{\text{loc}}(\xi, \theta). \quad (3.5)$$

There are two important observations to be made here. First, for each $1 \leq \ell \leq n$, there is only one factor in $z_\ell(\xi, \pm y)$ (3.4) that (up to sign) contains the sine of the squared absolute value (2.14) of one of the complex variables in its numerator:

- For $z_\ell(\xi, y)$, it is the factor corresponding to $m = \ell + p$, whose numerator is

$$\text{sgn}(M) \sin(|u_\ell|^2). \quad (3.6)$$

- For $z_\ell(\xi, -y)$, it is the factor with $m = \ell + n - p$, whose the numerator is either

$$\sin(\pi - \text{sgn}(M)|u_{\ell+n-p}|^2) = \text{sgn}(M) \sin(|u_{\ell+n-p}|^2), \quad \text{if } 1 \leq \ell \leq p, \quad (3.7)$$

or

$$\sin(\pi - \text{sgn}(M)|u_{\ell-p}|^2) = \text{sgn}(M) \sin(|u_{\ell-p}|^2), \quad \text{if } p < \ell \leq n. \quad (3.8)$$

Here we made use of $\xi_1 + \dots + \xi_n = \pi$, $\sin(\pi - \alpha) = \sin(\alpha)$ and $\sin(-\alpha) = -\sin(\alpha)$.

Second, the $(p-1)$ factors in $z_\ell(\xi, \pm y)$ with $m < \ell + p$ and $m > \ell + n - p$, respectively, are strictly negative and the factors corresponding to $m > \ell + p$ and $m < \ell + n - p$, respectively, are strictly positive for all ξ in the *closed* simplex \mathcal{A}_y . In particular, for any $\xi \in \mathcal{A}_y^+$ the sign of the ξ -dependent product in (3.4) equals $(-1)^{p-1} \text{sgn}(M) = \text{sgn}(\sin(ny))$, and therefore

$$z_\ell(\xi, \pm y) \geq 0, \quad \forall \xi \in \mathcal{A}_y, \quad \ell = 1, \dots, n. \quad (3.9)$$

We saw that z_ℓ can only vanish due to the numerators (3.6) and (3.7), (3.8), respectively. Consequently, in (3.4) the positive square root of $z_\ell(\xi, \pm y)$ can be taken for any $\xi \in \mathcal{A}_y^+$.

Now notice that, for all $\xi \in \mathcal{A}_y^+$, we have

$$v_j(\xi, y) = |u_j| w_j(\xi, y), \quad 1 \leq j \leq n, \quad (3.10)$$

where the $w_j(\xi, y)$ are positive and smooth functions of the form

$$w_j(\xi, y) = \left[\frac{\sin(|u_j|^2)}{|u_j|^2} \frac{(-1)^{p-1}}{\sin(\sum_{k=j}^{j+p-1} \xi_k)} \prod_{\substack{m=j+1 \\ (m \neq j+p)}}^{j+n-1} \frac{\sin(\sum_{k=j}^{m-1} \xi_k - y)}{\sin(\sum_{k=j}^{m-1} \xi_k)} \right]^{\frac{1}{2}}. \quad (3.11)$$

Similarly, we have

$$v_\ell(\xi, -y) = \begin{cases} |u_{\ell+n-p}| w_\ell(\xi, -y), & \text{if } 1 \leq \ell \leq p, \\ |u_{\ell-p}| w_\ell(\xi, -y), & \text{if } p < \ell \leq n \end{cases} \quad (3.12)$$

with the positive and smooth functions

$$w_\ell(\xi, -y) = \left[\frac{\sin(|u_{\ell+n-p}|^2)}{|u_{\ell+n-p}|^2} \frac{(-1)^{p-1}}{\sin(\sum_{k=\ell}^{\ell+n-p-1} \xi_k)} \prod_{\substack{m=\ell+1 \\ (m \neq \ell+n-p)}}^{\ell+n-1} \frac{\sin(\sum_{k=\ell}^{m-1} \xi_k + y)}{\sin(\sum_{k=\ell}^{m-1} \xi_k)} \right]^{\frac{1}{2}} \quad (3.13)$$

for $1 \leq \ell \leq p$, and

$$w_\ell(\xi, -y) = \left[\frac{\sin(|u_{\ell-p}|^2)}{|u_{\ell-p}|^2} \frac{(-1)^{p-1}}{\sin(\sum_{k=\ell}^{\ell+n-p-1} \xi_k)} \prod_{\substack{m=\ell+1 \\ (m \neq \ell+n-p)}}^{\ell+n-1} \frac{\sin(\sum_{k=\ell}^{m-1} \xi_k + y)}{\sin(\sum_{k=\ell}^{m-1} \xi_k)} \right]^{\frac{1}{2}} \quad (3.14)$$

for $p < \ell \leq n$.

The relation (2.17) allows us to express the ξ_k in terms of the complex variables for $k = 1, \dots, n-1$ as

$$\xi_k(u) = \sum_{j=1}^{n-1} \Omega_{j,k} (\text{sgn}(M) |u_j|^2 + c_j), \quad \text{with } c_j = \begin{cases} y, & \text{if } 1 \leq j \leq n-p, \\ y - \pi, & \text{if } n-p < j \leq n-1, \end{cases} \quad (3.15)$$

and $\xi_n(u) = \pi - \xi_1(u) - \dots - \xi_{n-1}(u)$. These formulas extend to $U(1)$ -invariant smooth functions on $S_{|M|}^{2n-1}$, which represent smooth functions on \mathbb{CP}^{n-1} on account of (2.29). By applying these, the above expressions $w_j(\xi(u), \pm y)$ ($j = 1, \dots, n$) give rise to smooth functions on \mathbb{CP}^{n-1} .

Definition 3.1. By setting $\theta_k = 0$ ($\forall k$) in the local Lax matrix L_y^{loc} (3.3) with y (2.8), we define the functions $\Lambda_{j,\ell}^y: \mathcal{A}_y^+ \rightarrow \mathbb{R}$ ($j, \ell = 1, \dots, n$) via the equations

$$\Lambda_{j,j+p}^y(\xi) = L_y^{\text{loc}}(\xi, \mathbf{1}_{n-1})_{j,j+p}, \quad 1 \leq j \leq n-p, \quad (3.16)$$

$$\Lambda_{j,j+p-n}^y(\xi) = L_y^{\text{loc}}(\xi, \mathbf{1}_{n-1})_{j,j+p-n}, \quad n-p < j \leq n, \quad (3.17)$$

$$\Lambda_{j,\ell}^y(\xi) = L_y^{\text{loc}}(\xi, \mathbf{1}_{n-1})_{j,\ell}(|u_j||u_{\ell+n-p}|)^{-1}, \quad 1 \leq j \leq n, \quad 1 \leq \ell \leq p \quad (\ell \neq j+p-n), \quad (3.18)$$

$$\Lambda_{j,\ell}^y(\xi) = L_y^{\text{loc}}(\xi, \mathbf{1}_{n-1})_{j,\ell}(|u_j||u_{\ell-p}|)^{-1}, \quad 1 \leq j \leq n, \quad p < \ell \leq n \quad (\ell \neq j+p). \quad (3.19)$$

The foregoing results lead to explicit formulas for $\Lambda_{j,\ell}^y$ (see Appendix A). Using the identification (2.29) and (3.15), it is readily seen that the $\Lambda_{j,\ell}^y(\xi(u))$ given by Definition 3.1 extend to smooth functions on \mathbb{CP}^{n-1} .

Remark 3.2. The explicit formulas of $\Lambda_{j,\ell}^y(\xi(u))$ contain products of square roots of strictly positive functions depending on $|u_k|^2 \in C^\infty(S_{|M|}^{2n-1})^{U(1)}$ for $k = 1, \dots, n$. In particular, they contain the square root of the function J given by

$$J(|u_k|^2) = \frac{\sin(|u_k|^2)}{|u_k|^2}, \quad (3.20)$$

which remains smooth (even real-analytic) at $|u_k|^2 = 0$ and is positive since we have $0 \leq |u_k|^2 \leq |M| < \pi$. Indeed, $|M| < \pi/q$ and $|M| < \pi/(n-q)$, respectively, for the two intervals of the type (i) couplings in (2.8).

The above observations allow us to introduce the following functions, which will be used to construct the global Lax matrix.

Definition 3.3. For $M > 0$ (2.10), define the smooth functions $L_{j,\ell}^{y,+}: \mathbb{CP}^{n-1} \rightarrow \mathbb{C}$ by

$$L_{j,\ell}^{y,+} \circ \pi_{|M|}(u) = \begin{cases} \Lambda_{j,j+p}^y(\xi(u)), & \text{if } 1 \leq j \leq n-p, \ell = j+p, \\ \Lambda_{j,j+p-n}^y(\xi(u)), & \text{if } n-p < j \leq n, \ell = j+p-n, \\ \bar{u}_j u_{\ell+n-p} \Lambda_{j,\ell}^y(\xi(u)), & \text{if } 1 \leq j \leq n, 1 \leq \ell \leq p, \ell \neq j+p-n, \\ \bar{u}_j u_{\ell-p} \Lambda_{j,\ell}^y(\xi(u)), & \text{if } 1 \leq j \leq n, p < \ell \leq n, \ell \neq j+p, \end{cases} \quad (3.21)$$

where u varies in $S_{|M|}^{2n-1}$. Then, for $M < 0$, define $L_{j,\ell}^{y,-}: \mathbb{CP}^{n-1} \rightarrow \mathbb{C}$ by

$$L_{j,\ell}^{y,-} \circ \pi_{|M|}(u) = L_{j,\ell}^{y,+} \circ \pi_{|M|}(\bar{u}), \quad (3.22)$$

referring to the right-hand-side of (3.21) with the understanding that now $y > p\pi/n$.

Next, we prove that the matrices L_y^{loc} and $L^{y,\pm} \circ \pi_{|M|} \circ \mathcal{E}$, are similar and can be transformed into each other by a unitary matrix. This is one of our main results.

Theorem 3.4. *The smooth matrix function $L^{y,\pm}: \mathbb{CP}^{n-1} \rightarrow \mathbb{C}^{n \times n}$ with components $L_{j,\ell}^{y,\pm}$ given by (3.21),(3.22) satisfies the following identity*

$$(L^{y,\pm} \circ \pi_{|M|} \circ \mathcal{E})(\xi, e^{i\theta}) = \Delta(e^{i\theta})^{-1} L_y^{\text{loc}}(\xi, e^{i\theta}) \Delta(e^{i\theta}), \quad \forall (\xi, e^{i\theta}) \in \mathcal{A}_y^+ \times \mathbb{T}^{n-1}, \quad (3.23)$$

where $\Delta(e^{i\theta}) = \text{diag}(\Delta_1, \dots, \Delta_n) \in U(n)$ with

$$\Delta_j = \exp\left(i \sum_{k=1}^{n-1} \Omega_{j,k} \theta_k\right), \quad j = 1, \dots, n-1, \quad \Delta_n = 1. \quad (3.24)$$

Consequently, $L^{y,\pm}(\pi_{|M|}(u)) \in \text{SU}(n)$ for every $u \in S_{|M|}^{2n-1}$, and $L^{y,\pm}$ provides an extension of the local Lax matrix L_y^{loc} (3.3) to the global phase space \mathbb{CP}^{n-1} .

Proof. The form of the local Lax matrix L_y^{loc} (3.3) and Definitions 3.1 and 3.3 show that (3.23) is equivalent to the equations

$$\Delta_j = \begin{cases} \Delta_{j+p}\rho_{j+p}, & \text{if } 1 \leq j \leq n-p, \\ \Delta_{j+p-n}\rho_{j+p-n}, & \text{if } n-p < j \leq n. \end{cases} \quad (3.25)$$

The two sides of (3.25) can be written as exponentials of linear combinations of the variables θ_k ($1 \leq k \leq n-1$). We next spell out the relations that ensure the exact matching of the coefficients of the θ_k in these exponentials. Plugging the components of Δ and ρ into (3.25), the case $1 \leq j < n-p$ gives

$$\begin{aligned} \Omega_{j,j+p-1} &= \Omega_{j+p,j+p-1} + 1, & (\text{coefficients of } \theta_{j+p-1}) \\ \Omega_{j,j+p} &= \Omega_{j+p,j+p} - 1, & (\text{coefficients of } \theta_{j+p}) \\ \Omega_{j,k} &= \Omega_{j+p,k}, & (\text{coefficients of } \theta_k, \ k \neq j+p-1, j+p), \end{aligned} \quad (3.26)$$

while for $j = n-p$ we get

$$\begin{aligned} \Omega_{n-p,n-1} &= 1, & (\text{coefficients of } \theta_{n-1}) \\ \Omega_{n-p,k} &= 0, & (\text{coefficients of } \theta_k, \ k \neq n-1). \end{aligned} \quad (3.27)$$

The case $n-p < j < n$ (and $p > 1$) leads to

$$\begin{aligned} \Omega_{j,j+p-n-1} &= \Omega_{j+p-n,j+p-n-1} + 1, & (\text{coefficients of } \theta_{j+p-n-1}) \\ \Omega_{j,j+p-n} &= \Omega_{j+p-n,j+p-n} - 1, & (\text{coefficients of } \theta_{j+p-n}) \\ \Omega_{j,k} &= \Omega_{j+p-n,k}, & (\text{coefficients of } \theta_k, \ k \neq j+p-n-1, j+p-n). \end{aligned} \quad (3.28)$$

For $j = n$ there are two possibilities. If $p = 1$ then we obtain

$$\begin{aligned} \Omega_{1,1} &= 1, & (\text{coefficients of } \theta_1) \\ \Omega_{1,k} &= 0, & (\text{coefficients of } \theta_k, \ k \neq 1), \end{aligned} \quad (3.29)$$

and if $p > 1$ then we require

$$\begin{aligned} \Omega_{p,p-1} &= -1, & (\text{coefficients of } \theta_{p-1}) \\ \Omega_{p,p} &= 1, & (\text{coefficients of } \theta_p) \\ \Omega_{p,k} &= 0, & (\text{coefficients of } \theta_k, \ k \neq p-1, p). \end{aligned} \quad (3.30)$$

Using the explicit formula given by Proposition 2.1, we now show that Ω satisfies (3.26). Since $\Omega_{j,k} = B_{j,k} - C_{j,k}$ for all j, k , where $B_{j,k}$ (2.21) depends on $(k-j)$ and $C_{j,k}$ (2.22) depends only on k , the equations (3.26) reduce to

$$\begin{aligned} B_{j,j+p-1} &= B_{j+p,j+p-1} + 1, \\ B_{j,j+p} &= B_{j+p,j+p} - 1, \\ B_{j,k} &= B_{j+p,k}, \quad k \neq j+p-1, j+p. \end{aligned} \quad (3.31)$$

The first equation holds, because $(j+p-1) - j = p-1 \equiv (n-q+1)p \pmod{n}$ implies $B_{j,j+p-1} = 1$ and $(j+p-1) - (j+p) = -1 \equiv (n-q)p \pmod{n}$ implies $B_{j+p,j+p-1} = 0$. For the second equation, we plainly have $B_{j,j+p} = 0$, and $(j+p) - (j+p) = 0 \equiv np \pmod{n}$ gives $B_{j+p,j+p} = 1$. Regarding the third equation, notice that $B_{j,k} = 0$ in (3.31) when $k-j \equiv \ell p \pmod{n}$ for some $\ell \in \{2, \dots, n-q\}$, and then $B_{j+p,k} = 0$ holds, too. Conversely,

$B_{j+p,k} = 0$ in (3.31) means that $(k-j) - p \equiv \ell p \pmod{n}$ for some $\ell \in \{1, \dots, n-q-1\}$, from which $(k-j) \equiv (\ell+1)p \pmod{n}$ and thus $B_{j,k} = 0$ follows. As B is a $(0,1)$ -matrix, we conclude that (3.31) is valid. Proceeding in a similar manner, we have verified the rest of the relations (3.27)–(3.30) as well. Since the relations (3.26)–(3.30) imply (3.25), the proof is complete. \square

It is an immediate consequence of Theorem 3.4 that the spectral invariants of the global Lax matrix $L^{y,\pm} \in C^\infty(\mathbb{CP}^{n-1}, \text{SU}(n))$ yield a Liouville integrable system. Because of (3.5) the corresponding Poisson commuting family contains the extension of the III_b Hamiltonian H_y^{loc} to \mathbb{CP}^{n-1} for any type (i) coupling. The self-duality of this compactified RS system was established in [8], and it will be studied in more detail elsewhere.

4 Compact forms of the elliptic RS system

In this section we explain that type (i) compactifications of the elliptic RS system can be constructed in exactly the same way as we saw for the trigonometric system. This is due to the fact that the local elliptic Lax matrix is built from the s-function (4.2) similarly as its trigonometric counterpart is built from the sine function, and on the real axis these two functions have the same zeros, signs, parity and antiperiodicity property.

We start by recalling some formulas of the relevant elliptic functions. First, let ω, ω' stand for the half-periods of the Weierstrass \wp function defined by

$$\wp(z; \omega, \omega') = \frac{1}{z^2} + \sum_{\substack{m,m'=-\infty \\ (m,m') \neq (0,0)}}^{\infty} \left[\frac{1}{(z - \omega_{m,m'})^2} - \frac{1}{\omega_{m,m'}^2} \right], \quad (4.1)$$

with $\omega_{m,m'} = 2m\omega + 2m'\omega'$. We adopt the convention $\omega, -i\omega' \in (0, \infty)$, which ensures that \wp is positive on the real axis. Next, introduce the following ‘s-function’:

$$s(z; \omega, \omega') = \frac{2\omega}{\pi} \sin\left(\frac{\pi z}{2\omega}\right) \prod_{m=1}^{\infty} \left[1 + \frac{\sin^2(\pi z/(2\omega))}{\sinh^2(m\pi|\omega'|/\omega)} \right], \quad (4.2)$$

related to the Weierstrass σ and ζ functions by $s(z) = \sigma(z) \exp(-\eta z^2/(2\omega))$ with the constant $\eta = \zeta(\omega)$. A useful identity connecting \wp and s is

$$\frac{s(z+z')s(z-z')}{s^2(z)s^2(z')} = \wp(z') - \wp(z), \quad z, z' \in \mathbb{C}. \quad (4.3)$$

The s-function is odd, has simple zeros at $\omega_{m,m'}$ ($m, m' \in \mathbb{Z}$) and enjoys the scaling property $s(tz; t\omega, t\omega') = t s(z; \omega, \omega')$. From now on we take

$$\omega = \frac{\pi}{2}, \quad (4.4)$$

whereby $s(z+\pi) = -s(z)$ holds as well. The trigonometric limit is obtained according to

$$\lim_{-i\omega' \rightarrow \infty} \wp(z; \pi/2, \omega') = \frac{1}{\sin^2(z)} - \frac{1}{3}, \quad \lim_{-i\omega' \rightarrow \infty} s(z; \pi/2, \omega') = \sin(z). \quad (4.5)$$

Let us now pick a type (i) coupling parameter y (2.8) and choose the domain of the dynamical variables to be the same $\mathcal{A}_y^+ \times \mathbb{T}^{n-1}$ as in the trigonometric case. Then consider the following IV_b variant of the standard [14, 17] elliptic RS Lax matrix:

$$L_y^{\text{loc}}(\xi, e^{i\theta}|\lambda)_{j,\ell} = \frac{s(y)s(x_j - x_\ell + \lambda)}{s(\lambda)s(x_j - x_\ell + y)} v_j(\xi, y) v_\ell(\xi, -y) \rho(\theta)_\ell, \quad \forall (\xi, e^{i\theta}) \in \mathcal{A}_y^+ \times \mathbb{T}^{n-1}, \quad (4.6)$$

where $\lambda \in \mathbb{C} \setminus \{\omega_{m,m'} : m, m' \in \mathbb{Z}\}$ is a spectral parameter and $v_\ell(\xi, \pm y) = \sqrt{z_\ell(\xi, \pm y)}$ with

$$z_\ell(\xi, \pm y) = \text{sgn}(s(ny)) \prod_{m=\ell+1}^{\ell+n-1} \frac{s(\sum_{k=\ell}^{m-1} \xi_k \mp y)}{s(\sum_{k=\ell}^{m-1} \xi_k)}. \quad (4.7)$$

These formulas are to be compared with the trigonometric case. Since $s(z)$ and $\sin(z)$ have matching properties on the real line, we can repeat the arguments presented in Section 3 to verify that $z_\ell(\xi, \pm y) > 0$ for every $\xi \in \mathcal{A}_y^+$. Taking positive square roots, and applying the relation $x_{k+1} - x_k = \xi_k$ to express $x_j - x_\ell$ in terms of ξ , we conclude that the above local Lax matrix is a smooth function on $\mathcal{A}_y^+ \times \mathbb{T}^{n-1}$ for every allowed value of the spectral parameter. The fact that it is a specialization of the standard elliptic Lax matrix ensures [14, 17] that its characteristic polynomial generates $(n-1)$ independent real Hamiltonians in involution with respect to the symplectic form (2.5). Indeed, the characteristic polynomial has the form

$$\det(L_y^{\text{loc}}(\xi, e^{i\theta}|\lambda) - \alpha \mathbf{1}_n) = \sum_{k=0}^n (-\alpha)^{n-k} c_k(\lambda, y) \mathcal{S}_k^{\text{loc}}(\xi, e^{i\theta}, y), \quad (4.8)$$

where the functions $\mathcal{S}_k^{\text{loc}}$ as well as their real and imaginary parts Poisson commute, and $\text{Re}(\mathcal{S}_k^{\text{loc}})$ for $k = 1, \dots, n-1$ are functionally independent. Explicit formulas of the c_k (that do not depend on the phase space variables) and $\mathcal{S}_k^{\text{loc}}$ (that do not depend on λ) can be found in [14, 17]. The function $\text{Re}(\mathcal{S}_1^{\text{loc}})$ is the RS Hamiltonian of IV_b type

$$\text{Re}(\text{tr} L_y^{\text{loc}}(\xi, e^{i\theta}|\lambda)) = \sum_{j=1}^n \cos(\theta_j - \theta_{j-1}) \sqrt{\prod_{m=j+1}^{j+n-1} \left[s(y)^2 (\wp(y) - \wp(\sum_{k=j}^{m-1} \xi_k)) \right]}. \quad (4.9)$$

We note in passing that in Ruijsenaars's papers [14, 17] one finds the elliptic Lax matrix $VL_y^{\text{loc}}V^{-1}$, where V is the diagonal matrix $V = \rho(\theta)\text{diag}(v_1(\xi, -y), \dots, v_n(\xi, -y))$. This difference is irrelevant, since it has no effect on the generated spectral invariants. Another difference is that we work in the center-of-mass frame.

Now the complete train of thought applied in the previous section remains valid if we simply replace the sine function with the s -function everywhere. In particular, the direct analogues of the formulas (3.10)–(3.14) hold with smooth functions $w_k(\xi, \pm y) > 0$, for $\xi \in \mathcal{A}_y$. Due to this fact, we can introduce a smooth elliptic Lax matrix defined on the global phase space \mathbb{CP}^{n-1} . The subsequent definition refers to the explicit formulas of Appendix A, which in the elliptic case contain the function

$$\mathcal{J}(|u_k|^2) = \frac{s(|u_k|^2)}{|u_k|^2}. \quad (4.10)$$

This has the same smoothness and positivity properties at and around zero as J (3.20) does. We also use $\xi(u)$ (3.15) and the functions $(x_j - x_\ell)(\xi)$ determined by $x_{k+1} - x_k = \xi_k$.

Definition 4.1. Take a type (i) y from (2.8) and represent the points of \mathbb{CP}^{n-1} as $\pi_{|M|}(u)$ with $u \in S_{|M|}^{2n-1}$. For $M > 0$ (2.10), define the smooth functions $\mathcal{L}_{j,\ell}^{y,+}$ on \mathbb{CP}^{n-1} by

$$\mathcal{L}_{j,\ell}^{y,+}(\pi_{|M|}(u)) = \begin{cases} \Lambda_{j,j+p}^y(\xi(u)), & \text{if } 1 \leq j \leq n-p, \ell = j+p, \\ \Lambda_{j,j+p-n}^y(\xi(u)), & \text{if } n-p < j \leq n, \ell = j+p-n, \\ \bar{u}_j u_{\ell+n-p} \Lambda_{j,\ell}^y(\xi(u)), & \text{if } 1 \leq j \leq n, 1 \leq \ell \leq p, \ell \neq j+p-n, \\ \bar{u}_j u_{\ell-p} \Lambda_{j,\ell}^y(\xi(u)), & \text{if } 1 \leq j \leq n, p < \ell \leq n, \ell \neq j+p, \end{cases} \quad (4.11)$$

with $\Lambda_{j,\ell}^y$ given in Appendix A. For $M < 0$, set $\mathcal{L}_{j,\ell}^{y,-}$ to be

$$\mathcal{L}_{j,\ell}^{y,-}(\pi_{|M|}(u)) = \mathcal{L}_{j,\ell}^{y,+}(\pi_{|M|}(\bar{u})) \quad (4.12)$$

with the understanding that in this case $y > p\pi/n$. Finally, define the λ -dependent elliptic Lax matrix $L^{y,\pm}$ on \mathbb{CP}^{n-1} by

$$L_{j,\ell}^{y,\pm}(\pi_{|M|}(u)|\lambda) = \frac{s((x_j - x_\ell)(\xi(u)) + \lambda)}{s(\lambda)} \mathcal{L}_{j,\ell}^{y,\pm}(\pi_{|M|}(u)), \quad (4.13)$$

where u runs over $S_{|M|}^{2n-1}$ and the spectral parameter λ varies in $\mathbb{C} \setminus \{\omega_{m,m'} : m, m' \in \mathbb{Z}\}$.

Theorem 4.2. *The spectral parameter dependent elliptic Lax matrix $L^{y,\pm}(\pi_{|M|}(u)|\lambda)$ (4.13) is a smooth global extension of $L_y^{\text{loc}}(\xi, e^{i\theta}|\lambda)$ (4.6) to the complex projective space \mathbb{CP}^{n-1} since it satisfies*

$$L^{y,\pm}((\pi_{|M|} \circ \mathcal{E})(\xi, e^{i\theta})|\lambda) = \Delta(e^{i\theta})^{-1} L_y^{\text{loc}}(\xi, e^{i\theta}|\lambda) \Delta(e^{i\theta}), \quad \forall (\xi, e^{i\theta}) \in \mathcal{A}_y^+ \times \mathbb{T}^{n-1}, \quad (4.14)$$

where Δ is given by (3.24) and $\pi_{|M|} \circ \mathcal{E} : \mathcal{A}_y^+ \times \mathbb{T}^{n-1} \rightarrow \mathbb{CP}^{n-1}$ is the symplectic embedding defined in Section 2.

The proof of Theorem 4.2 follows the lines of the proof of Theorem 3.4. The characteristic polynomial $\det(L^{y,\pm}(\pi_{|M|}(u)|\lambda) - \alpha \mathbf{1}_n)$ of the global Lax matrix depends smoothly on $\pi_{|M|}(u) \in \mathbb{CP}^{n-1}$ and as a consequence of (4.14) it satisfies

$$\det(L^{y,\pm}((\pi_{|M|} \circ \mathcal{E})(\xi, e^{i\theta})|\lambda) - \alpha \mathbf{1}_n) = \det(L_y^{\text{loc}}(\xi, e^{i\theta}|\lambda) - \alpha \mathbf{1}_n). \quad (4.15)$$

Since this holds for all α and λ , we see that the local IV_b Hamiltonian (4.9) together with its constants of motion $\text{Re}(\mathcal{S}_k^{\text{loc}})$, $k = 2, \dots, n-1$ extends to an integrable system on \mathbb{CP}^{n-1} . This was pointed out previously [17] for the special case $0 < y < \pi/n$ in (2.8).

In the trigonometric limit $-i\omega' \rightarrow \infty$ the s -function becomes the sine function, and we obtain a spectral parameter dependent trigonometric Lax matrix from the elliptic one. Then, setting the spectral parameter to be on the imaginary axis and taking the limit $-i\lambda \rightarrow \infty$ reproduces, up to conjugation by a diagonal matrix, the trigonometric global Lax matrix of Definition 3.3. Correspondingly, the global extension of the IV_b Hamiltonian (4.9) and its commuting family reduces to the global extension of the III_b Hamiltonian (2.6) and its constants of motion.

5 Conclusion and outlook

In this paper we have demonstrated by direct construction that the local phase space $\mathcal{A}_y^+ \times \mathbb{T}^{n-1}$ of the III_b and IV_b RS models (where \mathcal{A}_y^+ is the interior of the simplex (2.11)) can be embedded into \mathbb{CP}^{n-1} for any type (i) coupling y (2.8) in such a way that a suitable conjugate of the local Lax matrix extends to a smooth (actually real-analytic) function. Theorems 3.4 and 4.2 together with Appendix A provide explicit formulas for the resulting global Lax matrices. Their characteristic polynomials give rise to Poisson commuting real Hamiltonians on \mathbb{CP}^{n-1} that yield the Liouville integrable compactified trigonometric and elliptic RS systems.

Our direct construction was inspired by the earlier derivation of compactified III_b systems by quasi-Hamiltonian reduction [8]. The reduction identifies the III_b system with

a topological Chern-Simons field theory for any generic coupling parameter y . It appears natural to ask if an analogous derivation and relation to some topological field theory could exist for IV_b systems, too. We also would like to obtain a better understanding of the type (ii) trigonometric systems and their possible elliptic analogues.

In the near future, we wish to explore the classical dynamics and quantization of the III_b systems. This is partially motivated by the possibility to associate new random matrix ensembles with these systems [3]. For arbitrary type (i) couplings, geometric quantization yields the joint spectra of the quantized action variables effortlessly [7]. (It is necessary to introduce a second parameter into the systems before quantization, which can be achieved by taking an arbitrary multiple of the symplectic form.) The joint eigenfunctions of the quantized RS Hamiltonian and its commuting family should be derived by generalizing the results of van Diejen and Vinet [19].

Besides further studying the systems that we described, it would be also interesting to search for compactifications of generalized RS systems. We have in mind especially the BC_n systems due to van Diejen [18] and the recently introduced supersymmetric systems [2]. Regarding the former case, and even for general root systems, the results of [21] could be relevant, as well as the construction of Lax matrices for some of the BC_n systems reported in [12].

Throughout the text, we worked in the ‘center-of-mass frame’ and now we end by a comment on how the center-of-mass coordinate can be introduced into our systems. One possibility is to take the full phase space to be the Cartesian product of \mathbb{CP}^{n-1} with $\text{U}(1) \times \text{U}(1) = \{(e^{2iX}, e^{i\Phi})\}$ endowed with the symplectic form $|M|\omega_{\text{FS}} + dX \wedge d\Phi$. Here, e^{2iX} is interpreted as a center-of-mass variable for the n particles on the circle. Then n functions in involution result by adding an arbitrary function of $e^{i\Phi}$ to the $(n-1)$ commuting Hamiltonians generated by the ‘total Lax matrix’ $e^{-i\Phi} L^{y,\pm}$. On the dense open domain the total Lax matrix is obtained by replacing $\rho(\theta)$ in (4.6) by $\rho(\theta)e^{-i\Phi}$. By setting $e^{i\Phi}$ to 1 and quotienting by the canonical transformations generated by the functions of $e^{i\Phi}$ one recovers the phase space of the relative motion, \mathbb{CP}^{n-1} . There are also several other possibilities, as was discussed for analogous situations in [16, 5]. For example, one may replace $\text{U}(1) \times \text{U}(1)$ by its covering space $\mathbb{R} \times \mathbb{R}$.

Acknowledgements. We thank B.G. Puztai for helpful comments and S. Ruijsenaars for useful discussions. This work was supported in part by the Hungarian Scientific Research Fund (OTKA) under the grant K-111697 and by COST (European Cooperation in Science and Technology) in COST Action MP1405 QSPACE.

A Explicit form of the functions $\Lambda_{j,\ell}^y$

In this appendix we display the building blocks (4.11) of the global elliptic Lax matrix explicitly. Below, ξ varies in the closed simplex \mathcal{A}_y associated with a type (i) coupling y (2.8) for fixed p and M . The function \mathcal{J} was defined in (4.10). The trigonometric case is obtained by simply replacing the s-function (4.2) everywhere by the sine function.

Special components: For $1 \leq j \leq n-p$

$$\Lambda_{j,j+p}^y(\xi) = -\operatorname{sgn}(M) s(y) \frac{\left[\prod_{\substack{m=1 \\ (m \neq p)}}^{n-1} s(\sum_{k=j}^{j+m-1} \xi_k - y) s(\sum_{k=j+p}^{j+p+n-m-1} \xi_k + y) \right]^{\frac{1}{2}}}{\prod_{m=1}^{n-1} \left[s(\sum_{k=j}^{j+m-1} \xi_k) s(\sum_{k=j+p}^{j+p+m-1} \xi_k) \right]^{\frac{1}{2}}}.$$

For $n-p < j \leq n$

$$\Lambda_{j,j+p-n}^y(\xi) = \operatorname{sgn}(M) s(y) \frac{\left[\prod_{\substack{m=1 \\ (m \neq p)}}^{n-1} s(\sum_{k=j}^{j+m-1} \xi_k - y) s(\sum_{k=j+p-n}^{j+p-m-1} \xi_k + y) \right]^{\frac{1}{2}}}{\prod_{m=1}^{n-1} \left[s(\sum_{k=j}^{j+m-1} \xi_k) s(\sum_{k=j+p-n}^{j+p-m-1} \xi_k) \right]^{\frac{1}{2}}}.$$

Diagonal components: For $1 \leq j = \ell \leq p$

$$\Lambda_{j,j}^y(\xi) = [\mathcal{J}(|u_j|^2) \mathcal{J}(|u_{j+n-p}|^2)]^{\frac{1}{2}} \frac{\left[\prod_{\substack{m=1 \\ (m \neq p)}}^{n-1} s(\sum_{k=j}^{j+m-1} \xi_k - y) s(\sum_{k=j}^{j+n-m-1} \xi_k + y) \right]^{\frac{1}{2}}}{\prod_{m=1}^{n-1} s(\sum_{k=j}^{j+m-1} \xi_k)}.$$

For $p < j = \ell \leq n$

$$\Lambda_{j,j}^y(\xi) = [\mathcal{J}(|u_j|^2) \mathcal{J}(|u_{j-p}|^2)]^{\frac{1}{2}} \frac{\left[\prod_{\substack{m=1 \\ (m \neq p)}}^{n-1} s(\sum_{k=j}^{j+m-1} \xi_k - y) s(\sum_{k=j}^{j+n-m-1} \xi_k + y) \right]^{\frac{1}{2}}}{\prod_{m=1}^{n-1} s(\sum_{k=j}^{j+m-1} \xi_k)}.$$

Components above the diagonal: For $1 \leq j < \ell \leq p$

$$\Lambda_{j,\ell}^y(\xi) = s(y) [\mathcal{J}(|u_j|^2) \mathcal{J}(|u_{\ell+n-p}|^2)]^{\frac{1}{2}} \frac{\left[\prod_{\substack{m=1 \\ (m \neq \ell-j,p)}}^{n-1} s(\sum_{k=j}^{j+m-1} \xi_k - y) s(\sum_{k=\ell}^{\ell+n-m-1} \xi_k + y) \right]^{\frac{1}{2}}}{\prod_{m=1}^{n-1} \left[s(\sum_{k=j}^{j+m-1} \xi_k) s(\sum_{k=\ell}^{\ell+m-1} \xi_k) \right]^{\frac{1}{2}}}.$$

For $1 \leq j < \ell \leq n$ with $p < \ell$ and $\ell \neq j+p$

$$\Lambda_{j,\ell}^y(\xi) = \frac{s(y) [\mathcal{J}(|u_j|^2) \mathcal{J}(|u_{\ell-p}|^2)]^{\frac{1}{2}} \left[\prod_{\substack{m=1 \\ (m \neq \ell-j,p)}}^{n-1} s(\sum_{k=j}^{j+m-1} \xi_k - y) s(\sum_{k=\ell}^{\ell+n-m-1} \xi_k + y) \right]^{\frac{1}{2}}}{\operatorname{sgn}(j+p-\ell) \prod_{m=1}^{n-1} \left[s(\sum_{k=j}^{j+m-1} \xi_k) s(\sum_{k=\ell}^{\ell+m-1} \xi_k) \right]^{\frac{1}{2}}}.$$

Components below the diagonal: For $1 \leq \ell < j \leq n$ with $\ell \leq p$ and $\ell \neq j+p-n$

$$\Lambda_{j,\ell}^y(\xi) = \frac{s(y) [\mathcal{J}(|u_j|^2) \mathcal{J}(|u_{\ell+n-p}|^2)]^{\frac{1}{2}} \left[\prod_{\substack{m=1 \\ (m \neq j-\ell,p)}}^{n-1} s(\sum_{k=j}^{j+n-m-1} \xi_k - y) s(\sum_{k=\ell}^{\ell+m-1} \xi_k + y) \right]^{\frac{1}{2}}}{\operatorname{sgn}(\ell+n-j-p) \prod_{m=1}^{n-1} \left[s(\sum_{k=j}^{j+m-1} \xi_k) s(\sum_{k=\ell}^{\ell+m-1} \xi_k) \right]^{\frac{1}{2}}}.$$

For $p < \ell < j \leq n$

$$\Lambda_{j,\ell}^y(\xi) = s(y) [\mathcal{J}(|u_j|^2) \mathcal{J}(|u_{\ell-p}|^2)]^{\frac{1}{2}} \frac{\left[\prod_{\substack{m=1 \\ (m \neq j-\ell,p)}}^{n-1} s(\sum_{k=j}^{j+n-m-1} \xi_k - y) s(\sum_{k=\ell}^{\ell+m-1} \xi_k + y) \right]^{\frac{1}{2}}}{\prod_{m=1}^{n-1} \left[s(\sum_{k=j}^{j+m-1} \xi_k) s(\sum_{k=\ell}^{\ell+m-1} \xi_k) \right]^{\frac{1}{2}}}.$$

References

- [1] A. Alekseev, A. Malkin and E. Meinrenken, *Lie group valued moment maps*, J. Differ. Geom. **48**, 445-495 (1998); [arXiv:dg-ga/9707021](#)
- [2] O. Blondeau-Fournier, P. Desrosiers and P. Mathieu, *The supersymmetric Ruijsenaars-Schneider model*, Phys. Rev. Lett. **114**, 121602 (2015); [arXiv:1403.4667 \[hep-th\]](#)
- [3] E. Bogomolny, O. Giraud and C. Schmit, *Random matrix ensembles associated with Lax matrices*, Phys. Rev. Lett. **103**, 054103 (2009); [arXiv:0904.4898 \[nlin.CD\]](#)
- [4] P. Etingof, Calogero-Moser Systems and Representation Theory, European Mathematical Society, 2007
- [5] L. Fehér and V. Ayadi, *Trigonometric Sutherland systems and their Ruijsenaars duals from symplectic reduction*, J. Math. Phys. **51**, 103511 (2010); [arXiv:1005.4531 \[math-ph\]](#)
- [6] L. Fehér and C. Klimčík, *Self-duality of the compactified Ruijsenaars-Schneider system from quasi-Hamiltonian reduction*, Nucl. Phys. B **860**, 464-515 (2012); [arXiv:1101.1759 \[math-ph\]](#)
- [7] L. Fehér and C. Klimčík, *On the spectra of the quantized action-variables of the compactified Ruijsenaars-Schneider system*, Theor. Math. Phys. **171**, 704-714 (2012); [arXiv:1203.2864 \[math-ph\]](#)
- [8] L. Fehér and T.J. Kluck, *New compact forms of the trigonometric Ruijsenaars-Schneider system*, Nucl. Phys. B **882**, 97-127 (2014); [arXiv:1312.0400 \[math-ph\]](#)
- [9] A. Gorsky and N. Nekrasov, *Relativistic Calogero-Moser model as gauged WZW theory*, Nucl. Phys. B **436**, 582-608 (1995); [arXiv:hep-th/9401017](#)
- [10] N. Nekrasov, *Infinite-dimensional algebras, many-body systems and gauge theories*, pp. 263-299 in: Moscow Seminar in Mathematical Physics, AMS Transl. Ser. 2, eds. A.Yu. Morozov and M.A. Olshanetsky, American Mathematical Society, 1999
- [11] M.A. Olshanetsky and A.M. Perelomov, *Classical integrable finite-dimensional systems related to Lie algebras*, Phys. Rep. **71**, 313-400 (1981)
- [12] B.G. Puztai and T.F. Görbe, *Lax representation of the hyperbolic van Diejen dynamics with two coupling parameters*, preprint (2016); [arXiv:1603.06710 \[math-ph\]](#)
- [13] S.N.M. Ruijsenaars and H. Schneider, *A new class of integrable systems and its relation to solitons*, Ann. Phys. **170**, 370-405 (1986)
- [14] S.N.M. Ruijsenaars, *Complete integrability of relativistic Calogero-Moser systems and elliptic function identities*, Commun. Math. Phys. **110**, 191-213 (1987)
- [15] S.N.M. Ruijsenaars, *Finite-dimensional soliton systems*, pp. 165-206 in: Integrable and Superintegrable Systems, ed. B. Kupershmidt, World Scientific, 1990

- [16] S.N.M. Ruijsenaars, *Action-angle maps and scattering theory for some finite-dimensional integrable systems. III. Sutherland type systems and their duals*, Publ. RIMS Kyoto Univ. **31**, 247-353 (1995)
- [17] S.N.M. Ruijsenaars, *Systems of Calogero-Moser type*, pp. 251-352 in: Proceedings of the 1994 CRM-Banff Summer School ‘Particles and Fields’, eds. G.W. Semenoff and L. Vinet, Springer, 1999
- [18] J.F. van Diejen, *Deformations of Calogero-Moser systems*, Theor. Math. Phys. **99**, 549-554 (1994); [arXiv:solv-int/9310001](#)
- [19] J.F. van Diejen and L. Vinet, *The quantum dynamics of the compactified trigonometric Ruijsenaars-Schneider model*, Commun. Math. Phys. **197**, 33-74 (1998); [arXiv:math/9709221](#) [[math-ph](#)]
- [20] J.F. van Diejen and L. Vinet (eds), *Calogero-Moser-Sutherland Models*, Springer, 2000
- [21] J.F. van Diejen and E. Emsiz, *Orthogonality of Macdonald polynomials with unitary parameters*, Math. Z. **276**, 517-542 (2014); [arXiv:1301.1276](#) [[math.RT](#)]